

## Solutions to Chen's Plasma Physics

Kalman Krizhník

**1-1.** Compute the density (in units of  $m^{-3}$ ) of an ideal gas under the following conditions:

a) At  $0^\circ\text{C}$  and 760 Torr pressure (1 Torr = 1mm Hg). This is called the Loschmidt number.

b) In a vacuum at  $10^{-3}$  Torr at room temperature ( $20^\circ\text{C}$ ). This number is a useful one for the experimentalist to know by heart ( $10^{-3}$  Torr = 1 micron).

a) Avogadro's number is  $N_A = 6.022 \times 10^{23}$ . One mole of gas at STP occupies 22.4 liters. 1 liter is  $1 \times 10^{-3}$  cubic meters. Thus, the number per cubic meter is  $N_A/n = 6.022 \times 10^{23}/(22.4 \times 10^{-3}) = 2.66 \times 10^{25} m^{-3}$ . Thus, the Loschmidt number is  $2.66 \times 10^{25} \square$

b) Using  $PV=NkT$ , we obtain (with  $R=1.4 \times 10^{-23} J K^{-1}$  and 1 Torr = 133 Pa):

$$n = \frac{N}{V} = \frac{P}{kT} = \frac{10^{-3} \times 133}{1.4 \times 10^{-23} \times (20 + 273)} = 3.3 \times 10^{19} m^{-3} \quad \square \quad (1)$$

**1-2.** Derive the constant  $A$  for a normalized one-dimensional Maxwellian distribution

$$f(u) = Ae^{-mu^2/2kT} \quad (2)$$

such that

$$\int_{-\infty}^{\infty} f(u)du = 1 \quad (3)$$

This one is straightforward. Just integrate:

$$1 = \int_{-\infty}^{\infty} Ae^{-mu^2/2kT} du = A\sqrt{\frac{2\pi kT}{m}} \Rightarrow A = \sqrt{\frac{m}{2\pi kT}} \quad \square \quad (4)$$

**1-4.** Compute the pressure, in atmospheres and in  $tons/ft^2$ , exerted by a thermonuclear plasma on its container. Assume  $kT_e = kT_i = 20keV$ ,  $n = 10^{21}m^{-3}$  and  $p = nkT$ , where  $T = T_i + T_e$ .

This is just unit conversion, albeit with units that nobody really ever remembers. For reference,  $1 keV = 1.6 \times 10^{-19}J$ , so

$$p = 10^{21} \times (20keV + 20keV) = 4 \times 10^{22}m^{-3} keV = 4 \times 10^3m^{-3} J = 4 \times 10^5 N/m^2 \quad (5)$$

But  $1 atm = 10^5 N/m^2 = 1 ton/ft^2$ , so (Note: I think there is a mistake in Chen's solutions here. If I am mistaken, please let me know).

$$p = 0.04 atm = 0.04 ton/ft^2 \quad \square \quad (6)$$

**1-5.** In a strictly steady state situation, both the ions and the electrons will follow the Boltzmann relation

$$n_j = n_0 e^{-e_j\phi/kT_j} \quad (7)$$

For the case of an infinite, transparent grid charged to a potential  $\phi$ , show that the shielding distance is given approximately by

$$\lambda_D^{-2} = \frac{ne^2}{\epsilon_0} \left( \frac{1}{kT_e} + \frac{1}{kT_i} \right) \quad (8)$$

Show that  $\lambda_D$  is determined by the temperature of the colder species.

We'll use Poisson's equation

$$\nabla^2\phi = \frac{en_e - en_i}{\epsilon_0} = \frac{en_0}{\epsilon_0} (e^{e\phi/kT_e} - e^{-e\phi/kT_i}) \approx \frac{en_0}{\epsilon_0} \left( 1 + \frac{e\phi}{kT_e} - 1 - \frac{-e\phi}{kT_i} \right) = \frac{e^2n}{\epsilon_0} \left( \frac{\phi}{kT_e} + \frac{\phi}{kT_i} \right) \quad (9)$$

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